Universe in a helium droplet topological Fermi liquids: from Migdal jump to topological Khodel fermion condensate



Aalto University

G. Volovik

Landau Institute

C.D Landau



Migdal-100, Chernogolovka June 25, 2011

- * Topological media as topological objects in *p*-space
- * Fermi surface as vortex in *p*-space
- * Topological protection of Migdal jump and other singularities on Fermi surface
- * Khodel Fermi condensate as π -vortex
- Bulk-surface & bulk-vortex correspondence in topological matter
- * New life of Fermi condensate: Topologically protected flat bands
- * Can Fermi surface terminate? Fermi-arc in topological matter with Weyl points
- 3He Universe

classes of topological matter as momentum-space objects



Weyl point - hedgehog in **p**-space 3He-A, vacuum of SM, topological semimetals



flat band (Khodel state): π -vortex in **p**-space



Fermi arc on 3He-A surface & flat band on vortex: Dirac strings in **p**-space terminating on monopole





Migdal jump & p-space topology

* Singularity at Fermi surface is robust to perturbations: winding number N=1 cannot change continuously, interaction (perturbative) cannot destroy singularity

* Typical singularity: Migdal jump



* Other types of singularity: Luttinger Fermi liquid, marginal Fermi liquid, pseudo-gap ...

$$G(\omega, \mathbf{p}) = \frac{Z(p, \omega)}{i\omega - \varepsilon(p)}$$
$$Z(p, \omega) = (\omega^2 + \varepsilon^2(p))^{\gamma}$$

* Zeroes in Green's function instead of poles (for $\gamma > 1/2$) have the same winding number N=1



From Migdal jump to Landau Fermi-liquid



"Stability conditions & Fermi surface topologies in a superconductor" Gubankova-Schmitt-Wilczek, *Phys.Rev.* **B**74 (2006) 064505





non-topological flat bands due to interaction *Khodel-Shaginyan fermion condensate* JETP Lett. **51**, 553 (1990) GV, JETP Lett. 53, 222 (1991) Nozieres, J. Phys. (Fr.) 2, 443 (1992) $\delta E\{n(p)\} = \int \varepsilon(p) \delta n(p) d^{d}p = 0$ $E\{n(p)\}$ solutions: $\varepsilon(p) = 0$ or $\delta n(p) = 0$ $\epsilon(p)$ *n*(*p*) n(p)flat band p_F p_2 p_1 $\delta n(p) = 0$ $\delta n(p) = 0$ $\varepsilon(p) = 0$ $\delta n(p) = 0$ $\varepsilon(p) = 0 \quad \delta n(p) = 0$

splitting of Fermi surface to flat band

Flat band as momentum-space dark soliton terminated by half-quantum vortices



phase of Green's function changes by π across the "dark soliton"

topological correspondence:

topology in bulk protects gapless fermions on the surface or in vortex core

bulk-surface correspondence:

2D Quantum Hall insulator & 3He-A film

3D topological insulator

superfluid 3He-B

superfluid 3He-A, Weyl point semimetal

graphene

semimetal with Fermi lines

bulk-vortex correspondence:

superfluid 3He-A

chiral edge states

Dirac fermions on surface

Majorana fermions on surface

Fermi arc on surface

dispersionless 1D flat band on surface

2D flat band on the surface

1D flat band of zero modes in the core

Bulk-surface correspondence:

Edge states in quantum Hall topological insulators & 3He-A film



p-space skyrmion

p-space invariant in terms of Green's function & topological QPT



topology of graphene nodes

$$N = \frac{1}{4\pi i} \operatorname{tr} \left[\mathbf{K} \oint dl \, \mathbf{H}^{-1} \, \partial_l \, \mathbf{H} \right]$$

K - symmetry operator, commuting or anti-commuting with **H**

close to nodes:

$$\mathbf{H}_{N=+1} = \tau_{x}p_{x} + \tau_{y}p_{y}$$
$$\mathbf{H}_{N=-1} = \tau_{x}p_{x} - \tau_{y}p_{y}$$
$$\mathbf{K} = \tau_{z}$$



 E/γ_0

 k_x/a

exotic fermions: massless fermions with quardatic dispersion semi-Dirac fermions fermions with cubic and quartic dispersion

N = +1

bilayer graphene double cuprate layer surface of top. insulator neutrino physics

$$N = \frac{1}{4\pi i} \mathbf{tr} \left[\mathbf{K} \oint dl \mathbf{H}^{-1} \partial_l \mathbf{H} \right]$$







 $E^2 = 2c^2p^2 + 4m^2$



N=+2

 $E^2 = (p^2/2m)^2$

Dirac fermions

massive fermions

massless fermions with quadratic dispersion

multiple Fermi point

T. Heikkilä & GV arXiv:1010.0393



multilayered graphene

$$N = 1 + 1 + 1 + \dots$$

$$E = p^{N}$$
$$E = -p^{N}$$

spectrum in the outer layers

what kind of induced gravity emerges near degenerate Fermi point?

route to topological flat band on the surface of 3D material

Nodal spiral generates flat band on the surface

projection of spiral on the surface determines boundary of flat band



$$N_1 = \frac{1}{4\pi i} \operatorname{tr} \left[\mathbf{K} \oint_{\mathbf{C}} dl \, \mathbf{H}^{-1} \, \partial_l \, \mathbf{H} \right]$$

at each (p_x, p_y) except the boundary of circle one has 1D gapped state (insulator)

 $N_{\text{outside}} = 0$ trivial 1D insulator

 $N_{\text{inside}} = 1$ topological 1D insulator

at each (p_x, p_y) inside the circle one has 1D gapless edge state this is flat band

Nodal spiral generates flat band on the surface

projection of nodal spiral on the surface determines boundary of flat band

lowest energy states: surface states form the flat band

energy spectrum in bulk (projection to p_x , p_y plane)



Gapless topological matter with protected flat band on surface or in vortex core



non-topological flat bands due to interaction *Khodel-Shaginyan fermion condensate* JETP Lett. **51**, 553 (1990) GV, JETP Lett. **53**, 222 (1991) Nozieres, J. Phys. (Fr.) **2**, 443 (1992)



Bulk-surface correspondence:

Flat band on the surface of topological semimetals



flat band on the surface

projection of spiral on the surface determines boundary of the flat band

topologically protected nodal line in the form of spiral

$$N_1 = \frac{1}{4\pi i} \mathbf{tr} \left[\mathbf{K} \oint_{\mathbf{C}} dl \, \mathbf{H}^{-1} \, \partial_l \, \mathbf{H} \right]$$

flat band in soliton

$$H = \tau_3 (p_x^2 + p_z^2 - p_F^2)/2m + \tau_1 c(z)p_z$$

nodes at $p_z = 0$ and $p_x^2 = p_F^2$

$$N = \frac{1}{4\pi i} \mathbf{tr} \left[\mathbf{K} \oint dl \mathbf{H}^{-1} \partial_l \mathbf{H} \right]$$



Flat band on the graphene edge



Surface superconductivity in topological semimetals: route to room temperature superconductivity



Extremely high DOS of flat band gives high transition temperature:

normal superconductors: exponentially suppressed transition temperature

$$T_{c} = T_{F} \exp(-1/gv)$$

interaction DOS

$$1 = g \int \frac{d^2 p}{2h^2} \frac{1}{E(p)}$$

"Recent studies of the correlations between the internal microstructure of the samples and the transport properties suggest that superconductivity might be localized at the interfaces between crystalline graphite regions of different orientations, running parallel to the graphene planes." PRB. 78, 134516 (2008)

flat band superconductivity: linear dependence of T_c on coupling g



Stripes of increased diamagnetic susceptibility in underdoped superconducting $Ba(Fe_{1-x}Co_x)_2As_2$ single crystals: Evidence for an enhanced superfluid density at twin boundaries

B. Kalisky,^{1,2,*,†} J. R. Kirtley,^{1,2,3} J. G. Analytis,^{1,2,4} Jiun-Haw Chu,^{1,2,4} A. Vailionis,^{1,4} I. R. Fisher,^{1,2,4} and K. A. Moler^{1,2,4,5,*,‡}

Kathryn Moler: possible 2D superconductivityof twin boundaries



FIG. 1. (Color online) Local susceptibility image in underdoped Ba(Fe_{1-x}Co_x)₂As₂, indicating increased diamagnetic shielding on twin boundaries. (a) Local diamagnetic susceptibility, at T=17 K, of the *ab* face of sample UD1 (x=0.051 and $T_c=18.25$ K), showing stripes of enhanced diamagnetic response (white). In addition there is a mottled background associated with local T_c variations that becomes more pronounced as $T \rightarrow T_c$. Overlay: sketch of the scanning SQUID's sensor. The size of the pickup loop sets the spatial resolution of the susceptibility images. [(b) and (c)] Images of the same region at (b) T=17.5 K and (c) at T=18.5 K show that the stripes disappear above T_c . A topographic feature (scratch) appears in (b) and (c).

ဖွာ



relativistic quantum fields & gravity emerging near Weyl point

Atiyah-Bott-Shapiro construction:

linear expansion near Weyl point in terms of Dirac Γ -matrices

emergent relativity $H = e_a^k \Gamma^a \cdot (p_k - p_k^0)$

tetrad e_a^{μ}

secondary object: metric $g^{\mu\nu} = \eta^{ab} e_a^{\mu} e_b^{\nu}$

all ingredients of Standard Model : chiral fermions & gauge fields emerge in low-energy corner together with spin, Dirac Γ-matrices, gravity & physical laws: Lorentz & gauge invariance, equivalence principle, etc



From Weyl point to quantum Hall topological insulators



3D matter with Weyl points: Topologically protected flat band in vortex core

vortices in **r**-space



Topologically protected flat band in vortex core of superfluids with Weyl points



3He-A with Weyl points: Topologically protected Fermi arc on the surface

for each $|p_z| < p_F \cos \lambda$ one has 2D topological Hall insulator with zero energy edge states $E(p_z)=0$ (Dirac valley or Fermi arc PRB 094510,PRB 205101)



Fermi arc:

Fermi surface separates positive and negative energies, but has boundaries



Fermi surface of localized states is terminated by projections of Weyl points when localized states merge with

continuous spectrum

L spectrum of edge states on left wall



R spectrum of edge states on right wall

Conclusion: p-space topology determines fundamental many-body quantum phenomena

universality classes of quantum vacua

effective field theories in these quantum vacua (emergent gravity & QED)

topological quantum phase transitions (Lifshitz, plateau, etc.)

quantization of Hall and spin-Hall conductivity

topological Chern-Simons & Wess-Zumino terms

quantum statistics of topological objects

bulk-surface & bulk-vortex correspondence

exotic fermions: Majorana fermions; flat band; Fermi arc; cubic, quartic etc. dispersion

chiral anomaly & vortex dynamics, etc.

flat band & room-temperature superconductivity

superfuid phases ³He serve as primer for topological matter: quantum vacuum of Standard Model, topological superconductors, topological insulators, topological semimetals, etc.

we need: low T, high H, miniaturization, atomically smooth surface, nano-detectors, ... fabrication of samples for room-temperature superconductivity



